

SEMT30002 Scientific Computing and Optimisation

Week 3 Demos: 1D diffusion equations

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- In these demos, we showcase how the explicit and implicit Euler methods can be used to solve diffusion and reaction-diffusion equations.
- We start by importing some packages

```
In [1]: # importing packages
import numpy as np
import matplotlib.pyplot as plt
import matplotlib.animation as animation

# needed for animations in Jupyter notebook
%matplotlib notebook
```

Example 1 - explicit Euler

- We'll first look at how to use the explicit Euler method to solve the standard diffusion equation given by

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2}.$$

- The boundary and initial conditions are given by

$$u(a, t) = \alpha, \quad u(b, t) = \beta, \quad u(x, 0) = f(x).$$

- We will take $a = 0$ and $b = 1$; $\alpha = 1$ and $\beta = 0$; and $f(x) = 0$.

Approach

- The explicit Euler method will be implemented using nested `for` loops.
 - One loop is for time points, and the other loop is for grid points
- This approach does not take advantage of vectorisation so it is slow, but it's relatively straightforward to code up
- The first thing we do is define the problem parameters:

```
In [2]: """
        Define problem parameters
        """

        # Diffusion coefficient
        D = 0.5

        # Start and end of the spatial domain
        a = 0
        b = 1

        # Dirichlet boundary values
        alpha = 1.0
        beta = 0.0

        # Initial condition using lambda (anonymous) functions
        f = lambda x: np.zeros(np.size(x))
```

- The next thing we'll do is define the spatial grid.
- It is helpful to define the spatial grid first because this will determine the maximum time step that can be used with Euler's method

```
In [3]: # spatial grid
        N = 20                                # number of grid points (minus one)
        x = np.linspace(a, b, N+1)           # create the grid
        x_int = x[1:-1]                       # extract interior grid points
        dx = (b - a) / N                     # grid size
```

- We now compute the maximum size of the time step.
- This occurs when $C = D\Delta t/(\Delta x)^2 = 1/2$.
- So we first define $C_{\max} = 1/2$:

```
In [5]: # Numerical constant that determines stability
        C_max = 0.5
```

- From this value of C we can compute the maximum value that the time step can be:

```
In [6]: dt_max = C_max * dx**2 / D
        print(f'The largest time step can be {dt_max:.2e}')
```

The largest time step can be 2.50e-03

- We'll use a slightly smaller value of the time step to ensure stability.
- Moreover, we'll solve the problem until $t = 1$.
- We now discretise the time variable:

```
In [7]: # time discretisation
dt = 2.0e-3                # compute the time step
t_final = 1                # final time in the simulation
N_time = int(t_final / dt) # compute the number of time steps
t = dt * np.arange(N_time + 1) # compute time points

# Recalculate the constant C
C = D * dt / dx**2

# print some info about the time steps
print(N_time, 'time steps will be needed')
```

500 time steps will be needed

- Now that all of the variables have been defined, we can start solving the PDE using the explicit Euler method
- Recall that a `for` loop will be used to update the solution at each interior grid point
- We only loop over interior grid points because two Dirichlet boundary conditions are imposed

```
In [8]: """
        Time stepping using Euler's method
        """

# Pre-allocate the solution array. Rows are for space, columns for time
u = np.zeros((N - 1, N_time + 1))

# Set the first column of the solution array to the initial condition
u[:, 0] = f(x_int)

# loop over the time steps
for n in range(N_time):

    # loop over the grid
    for i in range(0, N-1):
        if i == 0:
            u[0, n+1] = u[0, n] + C * (u[1, n] - 2 * u[0, n] + alpha)
        elif 0 < i and i < N-2:
            u[i, n+1] = u[i, n] + C * (u[i+1, n] - 2 * u[i, n] + u[i-1, n])
        else:
            u[N-2, n+1] = u[N-2, n] + C * (beta - 2 * u[N-2, n] + u[N-3, n])
```

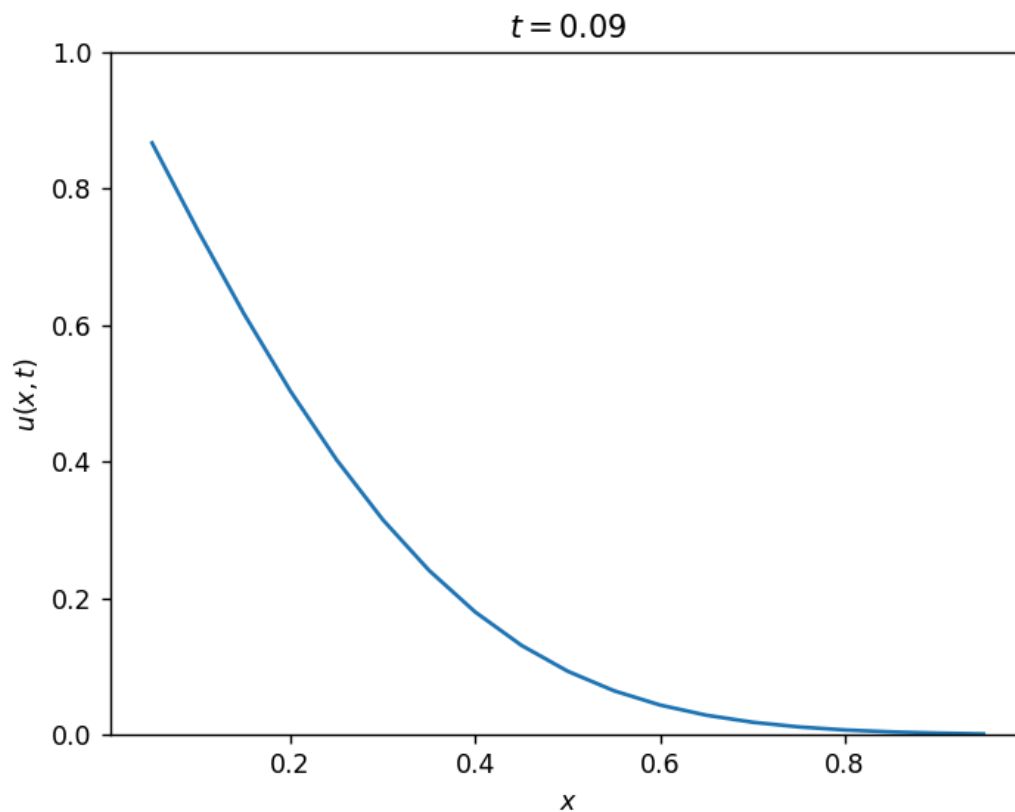
- We can use the `animation` module from `matplotlib` to create animations of the solution.
- We plot u as a function of space (x) and animate over time t
- If using VS Code with native Python files (.py), animations can be created without the `animation` module using the `plt.pause` and `plt.clf` functions from `Matplotlib`

```
In [9]: """
        animate the solution
        """
fig, ax = plt.subplots()
ax.set_ylim(0, 1)
ax.set_xlabel(f'$x$')
ax.set_ylabel(f'$u(x,t)$')

line, = ax.plot(x_int, u[:, 0])

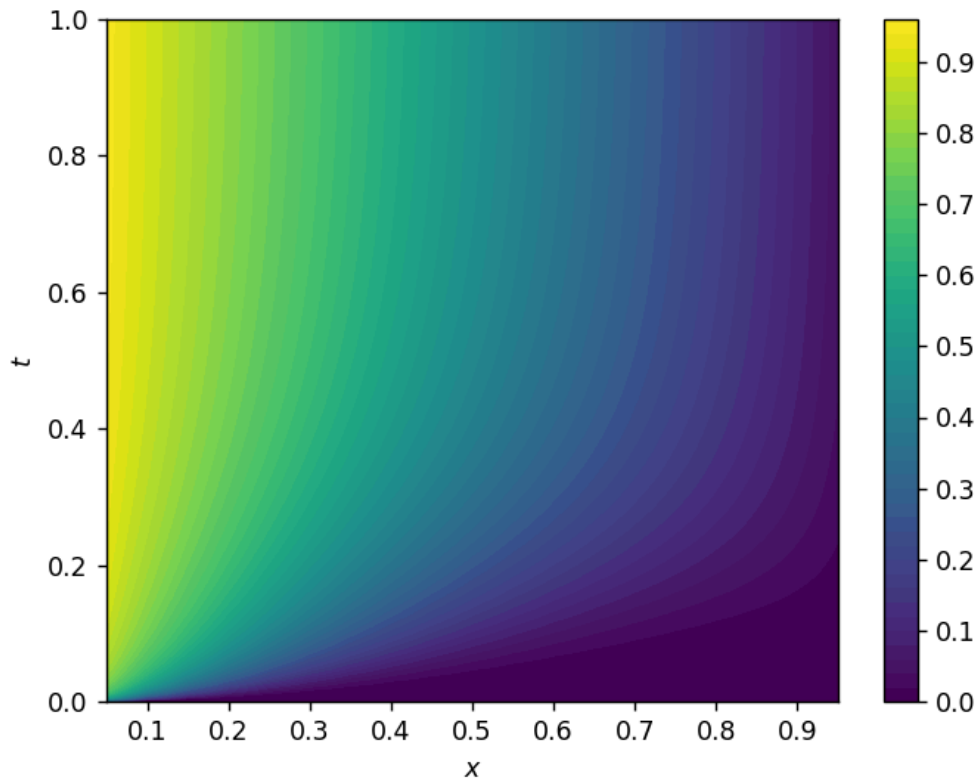
def animate(i):
    line.set_data((x_int, u[:, i]))
    ax.set_title(f'$t = {t[i]:.2f}$')
    return line

ani = animation.FuncAnimation(fig, animate, frames=N_time, blit=True, interval=100)
plt.show()
```



We can also create a contour plot to see the entire spatio-temporal evolution of the solution:

```
In [10]: plt.contourf(x_int, t, u.T, 50)
plt.xlabel('$x$')
plt.ylabel('$t$')
plt.colorbar()
plt.show()
```



Numerical blow up - choosing the wrong time step

- Although the explicit Euler method is simple, there is a strict restriction on the maximum size of the time step.
- This demo will now look at what happens when a time step that is too large is used with explicit Euler.
- Recall that the max time step is given by:

```
In [11]: C_max = 1/2
dt_max = C_max * dx**2 / D
print(f'The time step can be at most {dt_max:.2e}')
```

The time step can be at most 2.50e-03

- We'll now set the time step to be larger than this.
- Again, we'll attempt to solve the PDE until $t = 1$.

```
In [12]: # time discretisation
dt = 5.0e-3          # compute the time step
t_final = 1           # final time in the simulation
```

```

N_time = int(t_final / dt)      # compute the number of time steps
t = dt * np.arange(N_time + 1) # compute time points

# Recalculate the constant C
C = D * dt / dx**2

# print some info about the time steps
print(N_time, 'time steps will be needed')
print(f'The value of C is {C:.2f}')

```

200 time steps will be needed
The value of C is 1.00

- Now we'll re-solve the problem using explicit Euler:

```

In [13]: """
          Time stepping using Euler's method
          """

# Pre-allocate the solution array. Rows are for space, columns for time
u = np.zeros((N-1, N_time + 1))

# Set the first column of the solution array to the initial condition
u[:, 0] = f(x_int)

# loop over the time steps
for n in range(N_time):

    # loop over the grid
    for i in range(0, N-1):
        if i == 0:
            u[0, n+1] = u[0, n] + C * (u[1, n] - 2 * u[0, n] + alpha)
        elif 0 < i and i < N-2:
            u[i, n+1] = u[i, n] + C * (u[i+1, n] - 2 * u[i, n] + u[i-1, n])
        else:
            u[N-2, n+1] = u[N-2, n] + C * (beta - 2 * u[N-2, n] + u[N-3, n])

```

- Animating the solution shows that it becomes oscillatory, and the oscillation amplitude grows exponentially in time

```

In [14]: """
          animate the solution
          """

fig, ax = plt.subplots()
ax.set_ylim(-100, 100)
ax.set_xlabel(f'$x$')
ax.set_ylabel(f'$u(x,t)$')

line, = ax.plot(x_int, u[:, 0])

def animate(i):
    line.set_data((x_int, u[:, i]))
    ax.set_title(f'$t = {t[i]:.2f}$')

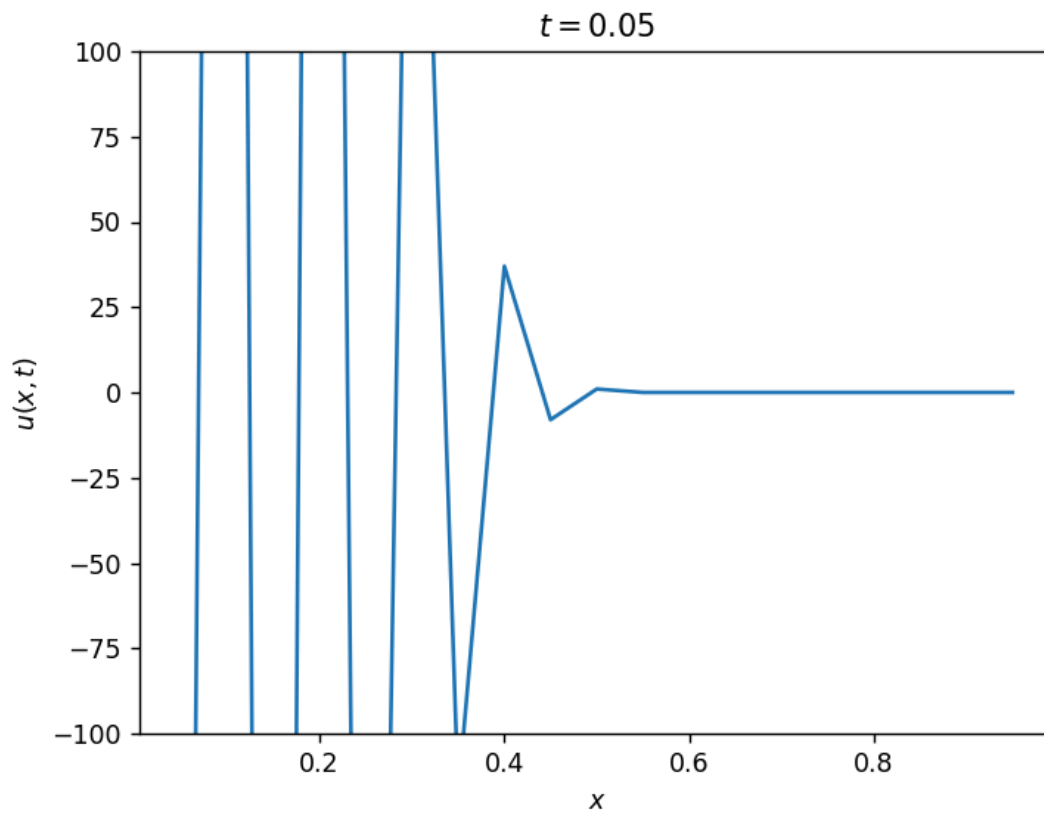
```

```

return line

ani = animation.FuncAnimation(fig, animate, frames=30, blit=True, interval=100)
plt.show()

```

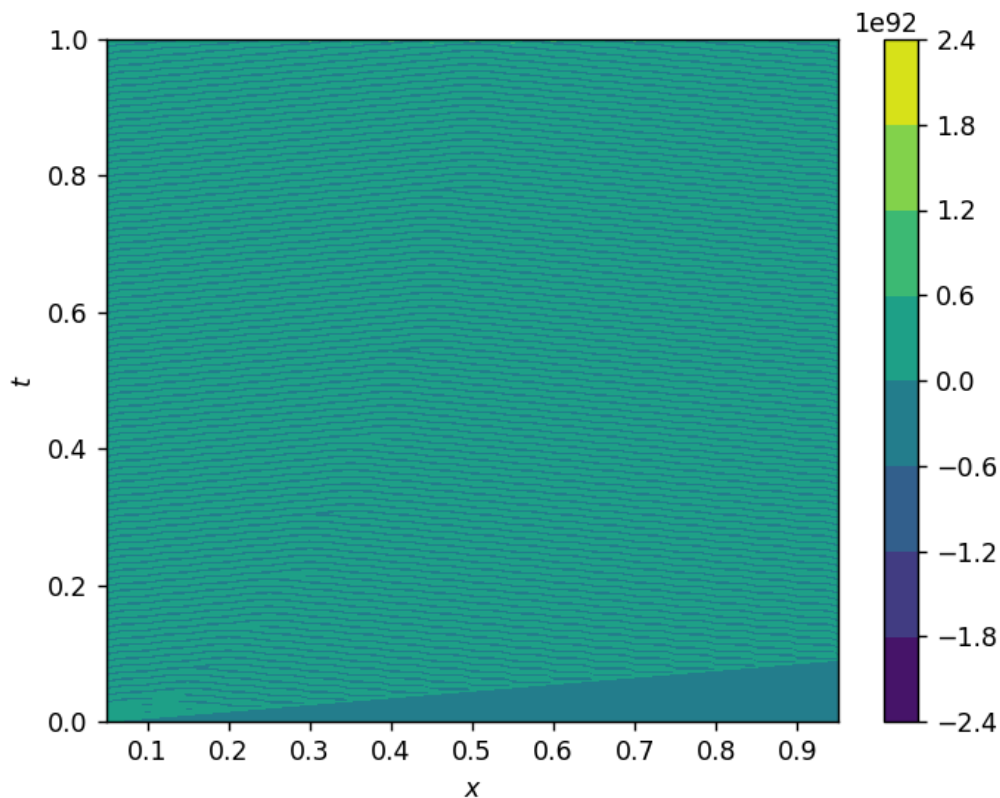


A contour plot also shows that the magnitude of the solution is becoming extremely large (this is why we say the solution is blowing up)

```

In [15]: plt.contourf(x_int, t, u.T)
plt.xlabel('$x$')
plt.ylabel('$t$')
plt.colorbar()
plt.show()

```



Example 2 - implicit Euler

- If we want to solve the diffusion equation with larger time steps, $\Delta t > (\Delta x)^2/2D$, then the implicit Euler method can be used.
- This involves solving a system of algebraic equations at each time step.

- In this demo we will solve the standard diffusion equation

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2}.$$

- The boundary and initial conditions will be

$$\left. \frac{\partial u}{\partial x} \right|_{x=0} = 0, \quad \left. \frac{\partial u}{\partial x} \right|_{x=1} = 0, \quad u(x, 0) = 1 + \cos(2\pi x)$$

- Unlike in other demos, two Neumann boundary conditions are being applied to this problem
- We first define the problem parameters and then the grid

```
In [16]: """
          Set up the problem parameters
```



```
"""
# Diffusion coefficient
D = 2

# Start and end of the spatial domain
a = 0
b = 1

# Initial condition using lambda (anonymous) functions
f = lambda x: 1 + np.cos(2 * np.pi * x)
```

```
In [17]: # spatial grid
N = 20                                     # number of grid points (minus one)
x = np.linspace(a, b, N+1)               # create the grid
dx = (b - a) / N                         # grid size
```

- Now we calculate the largest time step that would be possible using the **explicit** Euler method

```
In [18]: C_max = 1/2
dt_max = C_max * dx**2 / D
print(f'The time step can be at most {dt_max:.2e}')
```

The time step can be at most 6.25e-04

- We will use a time step that more than 10 times larger than this.
- This would lead to blow-up if the explicit Euler method was used
- However, the implicit Euler method will remain stable

```
In [19]: # time discretisation
dt = 1e-2                                # compute the time step
t_final = 1                              # final time in the simulation
N_time = int(t_final / dt)               # compute the number of time steps
t = dt * np.arange(N_time + 1)          # compute time points

# Recalculate the constant C
C = D * dt / dx**2

# print some info about the time steps
print(N_time, 'time steps will be needed')
print(f'The value of C is {C:.2f}')
```

100 time steps will be needed
The value of C is 8.00

- We will use NumPy's `linalg.solve` function to solve the linear algebraic system
- Now we generate the matrix \mathbf{A}^{NN}

```
In [20]: # Pre-allocate A^NN
A_NN = np.zeros((N+1, N+1))
```

```

# loop through each row
for n in range(N+1):

    # first row
    if n == 0:
        A_NN[n, n] = -2
        A_NN[n, n+1] = 2

    # last row
    elif n == N:
        A_NN[n, n-1] = 2
        A_NN[n, n] = -2

    # other rows
    else:
        A_NN[n, n-1] = 1
        A_NN[n, n] = -2
        A_NN[n, n+1] = 1

```

- We also generate the boundary condition vector \mathbf{b}^{NN} , and the identity matrix \mathbf{I} .
- The boundary conditions are of the form $\partial u / \partial x = 0$, so \mathbf{b}^{NN} is just an array of zeros.

```

In [21]: b_NN = np.zeros(N+1)
        I = np.eye(N+1)

```

We now start time stepping using the implicit Euler method

```

In [22]: """
        Time stepping using implicit Euler's method
        """

# Pre-allocate the solution array. Rows are for space, columns for time
u = np.zeros((N + 1, N_time + 1))

# Set the first column of the solution array to the initial condition
u[:, 0] = f(x)

# loop over the time steps
for n in range(N_time):

    # Solve the linear system
    u[:, n+1] = np.linalg.solve(I - C * A_NN, C * b_NN + u[:, n])

```

- As before, we animate the solution

```

In [23]: """
        animate the solution
        """

fig, ax = plt.subplots()
ax.set_ylim(0, 2)

```

```

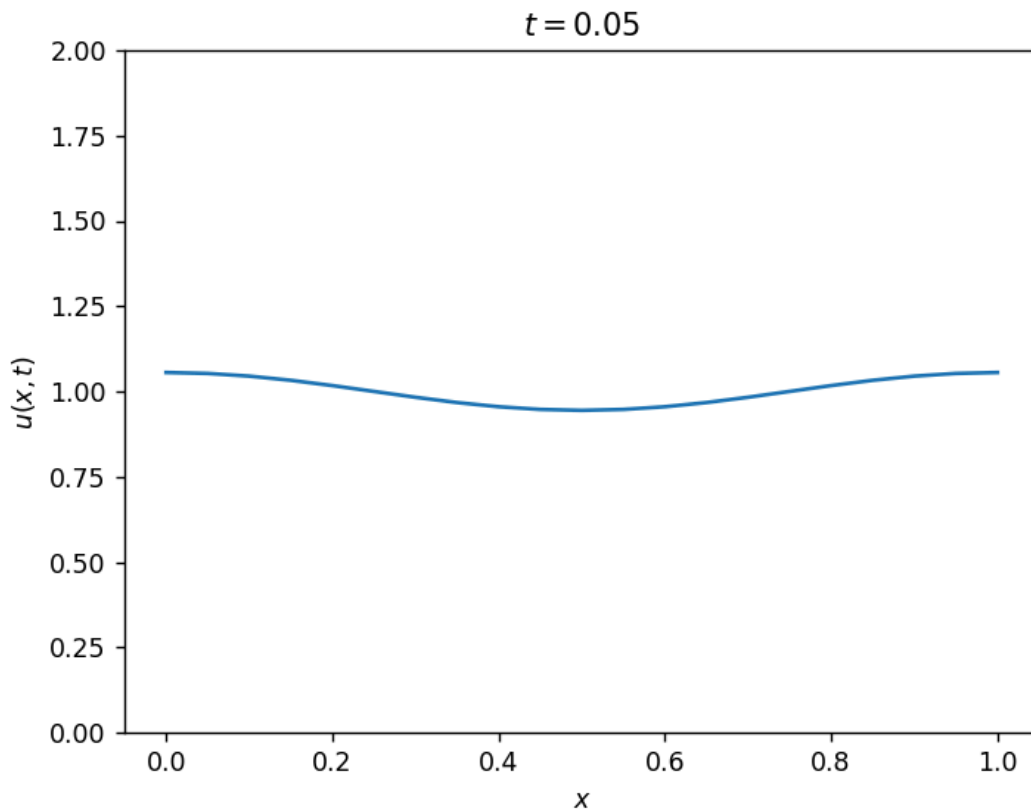
ax.set_xlabel(f'$x$')
ax.set_ylabel(f'$u(x,t)$')

line, = ax.plot(x, u[:, 0])

def animate(i):
    line.set_data((x, u[:, i]))
    ax.set_title(f'$t = {t[i]:.2f}$')
    return line

ani = animation.FuncAnimation(fig, animate, frames=30, blit=True, interval=100)
plt.show()

```



Validation

- For this problem it turns out that

$$\int_0^1 u(x, t) dx = 1$$

for all time.

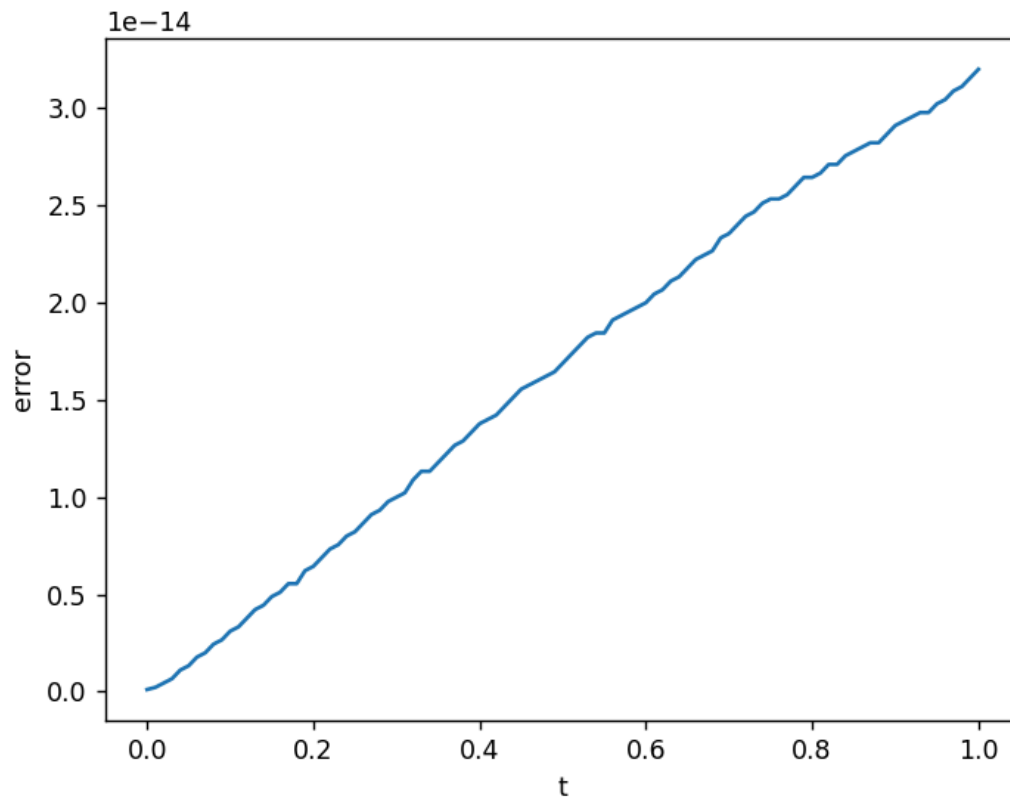
- In this case, the integral of u is called the mean of u
- We can check this is true using NumPy's `trapz` function as a means of validating the code

```
In [24]: # Pre-allocate space for the integral (mean) of u
mean_u = np.zeros(N_time+1)

# Loop over each time point
for n in range(N_time+1):

    # calculate the integral
    mean_u[n] = np.trapz(u[:, n], x)
```

```
In [25]: # Plot the results
plt.figure()
plt.plot(t, np.abs(mean_u - 1))
plt.xlabel('t')
plt.ylabel('error')
plt.show()
```



Example 3 - the Allen-Cahn equation

- The Allen-Cahn equation is a nonlinear reaction-diffusion equation given by

$$\frac{\partial u}{\partial t} = \varepsilon^2 \frac{\partial^2 u}{\partial x^2} + u - u^3$$

- Here, u is like a concentration, with $u = -1$ corresponding to pure water, $u = 1$ corresponding to pure oil, and $u = 0$ corresponding to a 50:50 mix of oil and water

- The initial condition will be $u(x, 0) = u_0(x)$, where $u_0(x) = 0 + \text{noise}$.
- The noise in the initial condition represents the fact there will be small fluctuations in the system.
- Noise is generated through an array of random numbers
- The boundary conditions are

$$\left. \frac{\partial u}{\partial x} \right|_{x=0,1} = 0.$$

Approach

- This PDE will be solved using the implicit Euler method
 - this requires solving a nonlinear algebraic system at each time step
 - the nonlinear system is solved using Newton's method
- This will be done using some modules I've created for solving PDEs
 - You don't have to do this, but I want to show you what is possible

```
In [26]: # import my modules for the unit
import finite_diff as fd
from pde_solvers import DiffusionEquation
```

- The boundary conditions are defined using a Python class:

```
In [27]: # Boundary conditions
bc_left = fd.BoundaryCondition('Neumann', lambda t: 0)
bc_right = fd.BoundaryCondition('Neumann', lambda t: 0)
```

- Then the space and time points are created using another class

```
In [28]: # Define the grid (space and time)
grid = fd.SpaceTimeGrid(N = 150, a = 0, b = 1, t_0 = 0, T = 50, dt = 1e-1)
```

- The parameters in the problem are defined
- The only parameter is ε^2 , which is equivalent to the diffusion coefficient

```
In [29]: # Define the parameter values
eps = 1e-2
pars = {"D": eps**2}
```

- Now Python functions that evaluate the initial condition and source terms are defined

```
In [30]: # Initial condition
def u_0(x, pars):
    return 1e-3 * (2 * np.random.random(grid.N+1) - 1)

# Source term
def q(u, x, pars):
    return u - u**3

# Derivative of source term wrt u (needed for Newton's method)
def q_prime(u, x, pars):
    return 1 - 3 * u**2
```

- Now the discretised equation is constructed using a Python class
 - The `DiffusionEquation` class generates the algebraic system for diffusion equations
 - All of the data structures are then stored in an object
- The algebraic system is then solved by calling the `solve` method

```
In [31]: # Define the PDE and solve
pde = DiffusionEquation(grid, bc_left, bc_right, u_0, pars, q, q_prime)
u = pde.solve(method = "implicit_Euler")
```

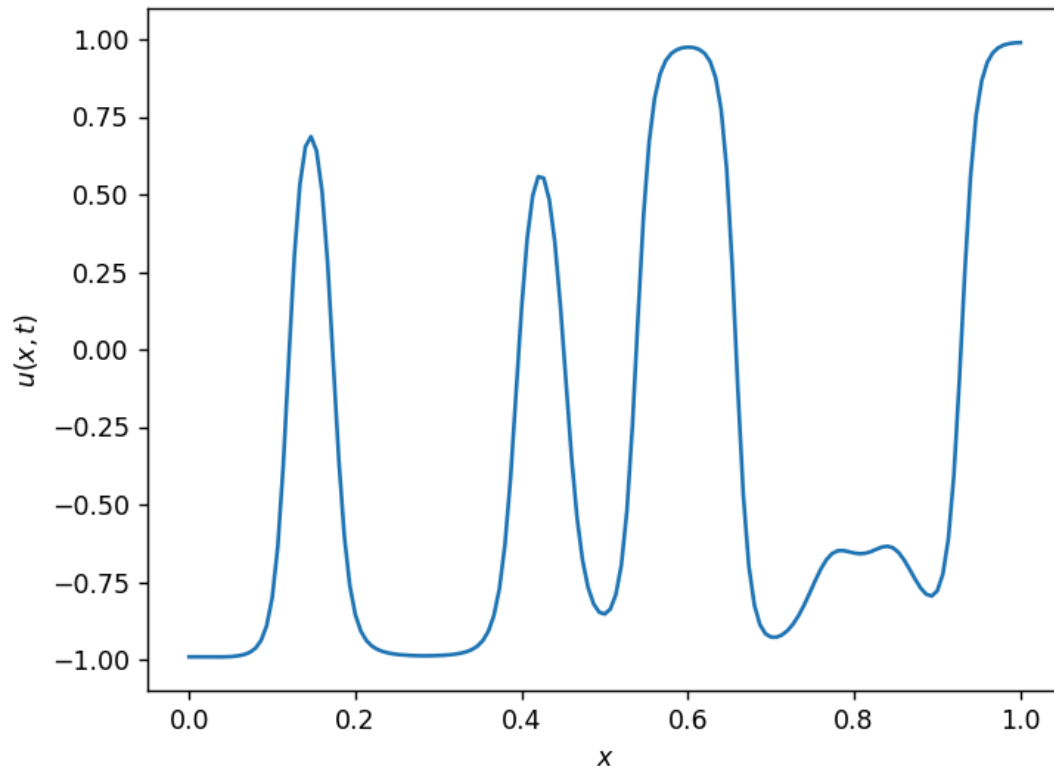
- Now the solution is visualised

```
In [32]: # Animate the solution
fig, ax = plt.subplots()
ax.set_ylim(-1.1, 1.1)
ax.set_xlabel('$x$')
ax.set_ylabel('$u(x,t)$')

line, = ax.plot(grid.x, u[:, 0])

def animate(i):
    line.set_data((grid.x, u[:, i]))
    return line

ani = animation.FuncAnimation(fig, animate, frames=grid.Nt, blit=True, interval=100)
plt.show()
```



```
In [33]: plt.figure()
plt.contourf(grid.x, grid.t, u.T, 50)
plt.xlabel('$x$')
plt.ylabel('$t$')
plt.colorbar()
plt.show()
```

